

Notes on DPLL*

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*Most of material is derived from Chistian's report[4]

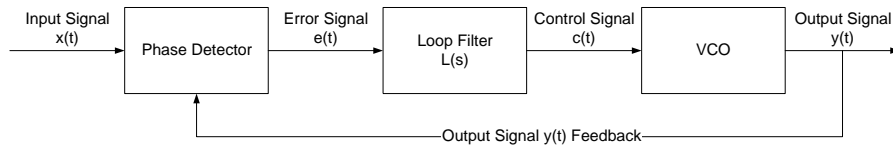


Figure 1: Structure of a PLL.

1 Digital Phase Locked Loop (DPLL)

1.1 Continuous-time Phase Locked Loop

Figure 1 shows the general structure of a Phase Locked Loop (PLL)[1]. A generic PLL consists of a phase detector, a loop filter, and a voltage-controlled oscillator (VCO). The phase detector produces the error signal that measures the difference between the input and output of the PLL. The loop filter averages the error and generates the control signal to the VCO to produce the output signal.

Usually the input and output signals of the PLL are sinusoids with input

$$x(t) = A_x \cos[2\pi f_c t + \phi(t)]$$

and output

$$y(t) = A_y \sin \left[2\pi f_c t + \hat{\phi}(t) \right]$$

After an ideal phase detector, the output error signal is

$$e(t) = \phi(t) - \hat{\phi}(t)$$

The loop filter is characterized in s -domain (Laplace transform) as

$$C(s) = L(s)E(s)$$

where $E(s)$ corresponds to the error signal $e(t)$, $L(s)$ is the transfer function of the loop filter, and $C(s)$ corresponds to the control signal $c(t)$. The VCO is basically a sinusoidal signal generator with an instantaneous carrier phase given by

$$2\pi f_c t + \hat{\phi}(t) = 2\pi f_c t + \int_{-\infty}^t c(\tau) d\tau$$

Notice we assume the VCO gain is unity here. Any VCO gain will be lumped into the loop filter gain. Taking the derivative of the above we have

$$\frac{d\hat{\phi}(t)}{dt} = c(t)$$

or equivalently in s -domain

$$s\hat{\Phi}(s) = C(s)$$

where on the LHS, we assumed $\hat{\phi}(0) = 0$. The loop transfer function is then

$$H(s) = \frac{\hat{\Phi}(s)}{\Phi(s)} = \frac{(1/s)C(s)}{\Phi(s)} = \frac{(1/s)L(s) [\Phi(s) - \hat{\Phi}(s)]}{\Phi(s)} = \frac{L(s)}{s} [1 - H(s)]$$

Solving the above, we have

$$H(s) = \frac{L(s)/s}{1 + L(s)/s} = \frac{L(s)}{s + L(s)}$$

A first-order PLL is characterized by the loop filter of the form

$$L(s) = K$$

and the transfer function becomes

$$H(s) = \frac{K}{s + K}$$

If we choose the loop filter of the form

$$L(s) = K \frac{s + K_1}{s + K_2}$$

the transfer function becomes

$$H(s) = \frac{Ks + KK_1}{s^2 + (K + K_2)s + KK_1}$$

which is a second-order PLL. It is customary to express the denominator of $H(s)$ in the form[2]

$$s^2 + 2\eta\omega_n s + \omega_n^2$$

where η is called the *loop damping factor* and ω_n is called the natural frequency of the loop. It is clear that

$$\omega_n = \sqrt{KK_1}$$

$$\eta = \frac{K + K_2}{2\sqrt{KK_1}}$$

the loop transfer function becomes

$$H(s) = \frac{Ks + \omega_n^2}{s^2 + 2\eta\omega_n s + \omega_n^2}$$

When $K_2 = 0$, we have a special second-order PLL with the loop filter

$$L(s) = K \frac{s + K_1}{s} \tag{1}$$

and transfer function

$$H(s) = \frac{Ks + KK_1}{s^2 + Ks + KK_1} = \frac{2\eta\omega_n s + \omega_n^2}{s^2 + 2\eta\omega_n s + \omega_n^2} \tag{2}$$

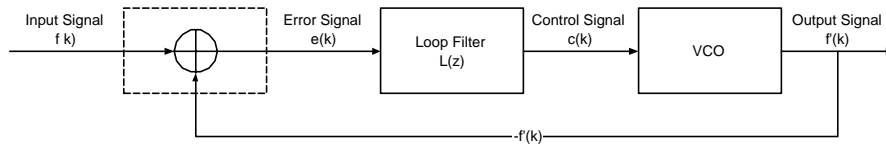


Figure 2: General structure of a DPLL.

and the natural frequency and damping factor are

$$\omega_n = \sqrt{KK_1}$$

$$\eta = \frac{1}{2}\sqrt{\frac{K}{K_1}}$$

respectively.

It is often useful to know the steady-state operating point of PLL. The steady-state phase error is defined to be

$$e_{ss} = \lim_{t \rightarrow \infty} e(t)$$

Using the final value theorem of Laplace transform, we have

$$e_{ss} = \lim_{s \rightarrow 0} sE(s) = \lim_{s \rightarrow 0} s\Phi(s) [1 - H(s)] = \lim_{s \rightarrow 0} \frac{s^2\Phi(s)}{s + L(s)}$$

1.2 General structure of DPLL

Figure 2 shows the structure of a Digital Phase Locked Loop. The phase detector is a simple subtractor. The error signal is

$$e(k) = \phi(k) - \hat{\phi}(k)$$

or in z -domain

$$E(z) = \Phi(z) - \hat{\Phi}(z)$$

Analogous to the differential equation for the VCO in the continuous PLL, we write the difference equation

$$\hat{\phi}(k+1) - \hat{\phi}(k) = c(k)$$

for the VCO in the DPLL. In z -domain, the above becomes

$$(z-1)\hat{\Phi}(z) = C(z) = L(z)E(z) = L(z) [\Phi(z) - \hat{\Phi}(z)]$$

which can be reorganized to give the transfer function

$$H(z) = \frac{L(z)}{L(z) + z - 1}$$

The steady state error is calculated as

$$e_{ss} = \lim_{k \rightarrow \infty} e(k) = \lim_{z \rightarrow 1} (z-1)E(z) = \lim_{z \rightarrow 1} \frac{(z-1)^2\Phi(z)}{L(z) + z - 1} \quad (3)$$

1.3 First-order DPLL

For the first-order DPLL,

$$L(z) = K$$

the transfer function becomes

$$H(z) = \frac{K}{K + z - 1}$$

The first-order DPLL is not quite suited for correcting the phase error caused by residual frequency offset as discussed in Section ??.

For example, if we assume the input sequence is of the form

$$\phi(n) = \epsilon_0 n$$

i.e. a residual frequency offset causing linear phase increment. To find out the z -transform of the sequence, we notice that the above sequence is expressed by the difference equation

$$\phi(k+1) - \phi(k) = \epsilon_0 \quad k \geq 0$$

and

$$\phi(k) = 0 \quad k \leq 0$$

The z -transform of the RHS of the difference equation is

$$\epsilon_0 \sum_{n=0}^{\infty} \frac{1}{z^n} = \epsilon_0 \frac{z}{z-1}$$

Performing z -transform on both sides of the difference equation yields

$$z\Phi(z) - \Phi(z) = \epsilon_0 \frac{z}{z-1}$$

and thus the z -transform of the input sequence $\phi(n)$ is

$$\Phi(z) = \epsilon_0 \frac{z}{(z-1)^2}$$

Substituting the sequence z -domain expression into the DPLL steady state error expression (3), we have

$$e_{ss} = \lim_{z \rightarrow 1} \frac{(z-1)^2 \epsilon_0 \frac{z}{(z-1)^2}}{K + z - 1} = \frac{\epsilon_0}{K}$$

The steady state error is non-zero but it can be reduced by increasing the loop gain K .

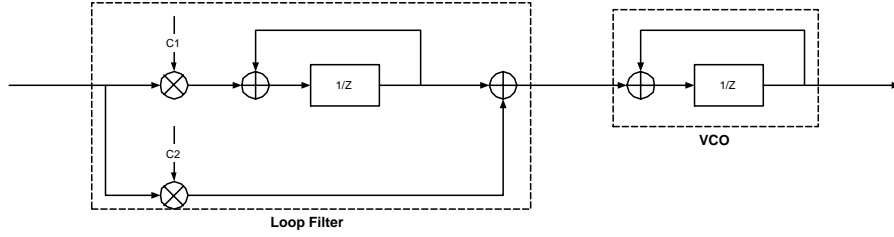


Figure 3: Second-order DPLL loop filter and VCO.

1.4 Second-order DPLL

1.4.1 Second-order DPLL structure

Noticing the discrete domain equivalent of s is $z-1$, we can construct a special second-order DPLL similar to the PLL discussed before, whose loop filter is, similar to (1),

$$L(z) = K \frac{(z-1) + K_1}{(z-1)}$$

and whose transfer function is, similar to (2),

$$H(z) = \frac{K(z-1) + KK_1}{(z-1)^2 + K(z-1) + KK_1}$$

The filter parameters can be matched with those used in [3] through

$$\begin{cases} K = C_2 \\ K_1 = C_1/C_2 \end{cases} \quad \text{i.e.} \quad \begin{cases} C_1 = KK_1 \\ C_2 = K \end{cases}$$

and the transfer function becomes

$$H(z) = \frac{C_2(z-1) + C_1}{(z-1)^2 + C_2(z-1) + C_1} \quad (4)$$

Figure 3 shows the block diagram of the second-order DPLL.

Now if the input is

$$\phi(n) = \epsilon_0 n$$

the steady state error is

$$e_{ss} = \lim_{z \rightarrow 1} \frac{(z-1)^2 \epsilon_0 \frac{z}{(z-1)^2}}{K \left(1 + \frac{1}{z-1} K_1\right) + (z-1)} = \lim_{z \rightarrow 1} \frac{\epsilon_0 z (z-1)}{(z-1)^2 + K(z-1) + KK_1} = 0$$

To express the DPLL in term natural frequency ω_n and damping factor η , we define

$$\begin{cases} C_1 = \omega_n^2 \\ C_2 = 2\eta\omega_n \end{cases} \quad \text{i.e.} \quad \begin{cases} \omega_n = \sqrt{C_1} \\ \eta = \frac{C_2}{2\sqrt{C_1}} \end{cases} \quad (5)$$

Notice here we require C_1 to be positive which as discussed later is a prerequisite for a stable DPLL. With the definition, we can express the DPLL transfer function as

$$H(z) = \frac{2\eta\omega_n(z-1) + \omega_n^2}{(z-1)^2 + 2\eta\omega_n(z-1) + \omega_n^2}$$

1.4.2 Second-order DPLL stable condition

To have a stable DPLL, all its poles must reside in the unit circle. The poles can be found by solving the equation in the denominator of $H(z)$,

$$(z-1)^2 + C_2(z-1) + C_1 = 0$$

yielding the solution of the form

$$z = 1 + \frac{-C_2 \pm \sqrt{C_2^2 - 4C_1}}{2} = \frac{2 - C_2 \pm \sqrt{C_2^2 - 4C_1}}{2}$$

When

$$\frac{C_2^2}{4} > C_1$$

the solution is real and we require

$$-1 < \frac{2 - C_2 \pm \sqrt{C_2^2 - 4C_1}}{2} < 1$$

Noticing that we only need to make the bigger pole to be less 1 and the smaller pole to be bigger than -1 , the above inequality is reduced to

$$\begin{aligned} \frac{2 - C_2 + \sqrt{C_2^2 - 4C_1}}{2} &< 1 \\ \frac{2 - C_2 - \sqrt{C_2^2 - 4C_1}}{2} &> -1 \end{aligned}$$

From the first of the two inequalities, we find

$$0 < \sqrt{C_2^2 - 4C_1} < C_2$$

and further

$$C_1 > 0$$

From the second of the two inequalities, we find

$$4 - C_2 > \sqrt{C_2^2 - 4C_1} > 0 \Rightarrow C_2 < 4$$

and further

$$C_1 > 2C_2 - 4$$

Combine all above we have

$$\begin{aligned} \frac{C_2^2}{4} &> C_1 \\ 0 &< C_2 < 4 \\ C_1 &> 0 \\ C_1 &> 2C_2 - 4 \end{aligned}$$

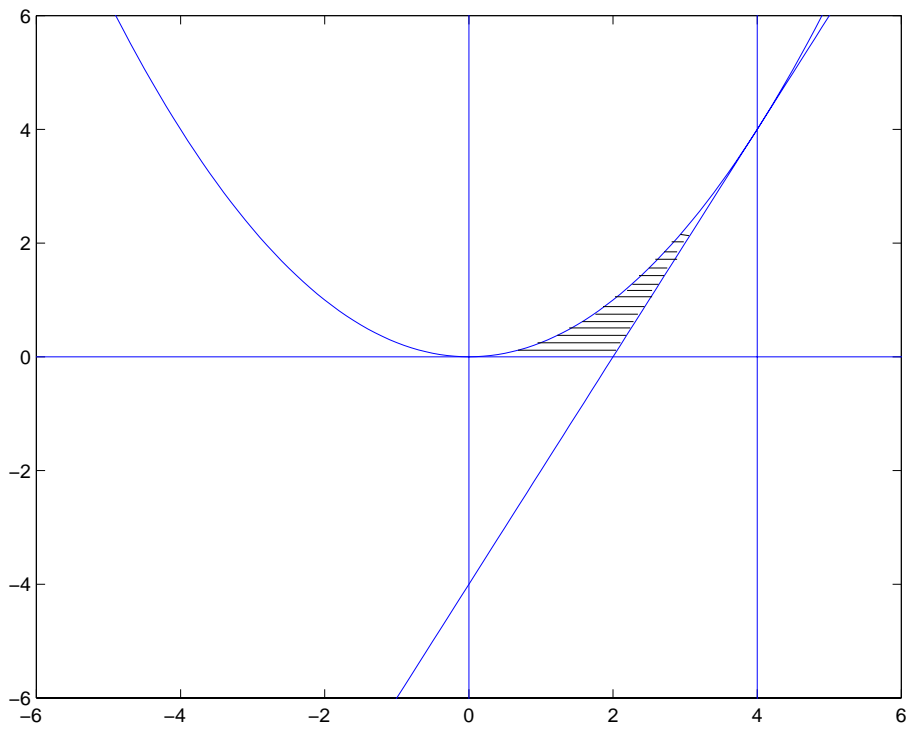


Figure 4: Second-order DPLL real solution region. Horizontal axis is C_2 . Vertical axis is C_1 .

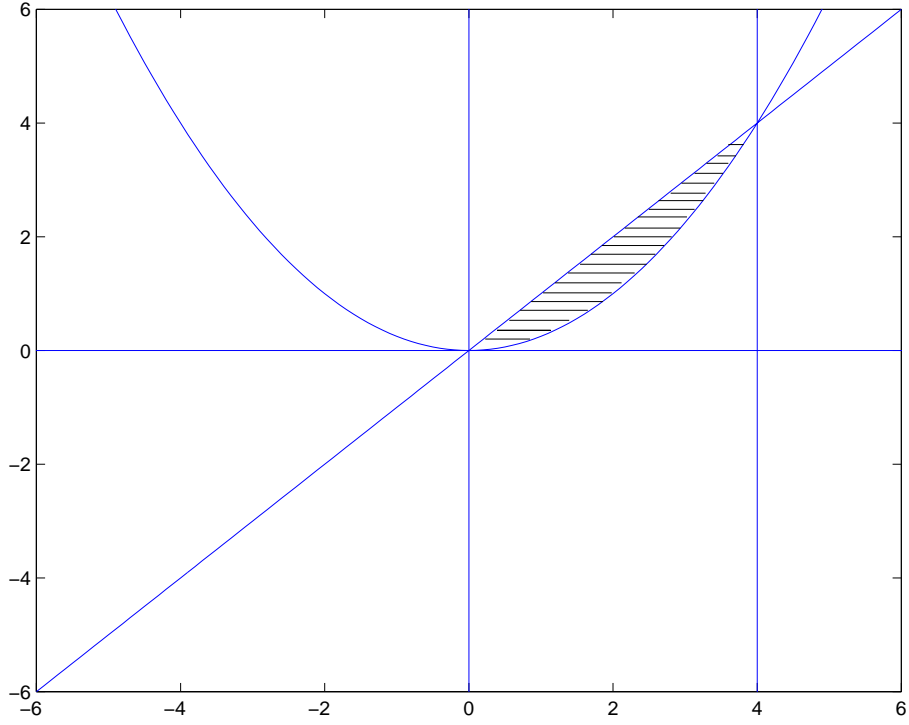


Figure 5: Second-order DPLL imaginary solution region. Horizontal axis is C_2 . Vertical axis is C_1 .

The shaded region in Figure 4 shows the solution region.

When

$$\frac{C_2^2}{4} \leq C_1$$

the poles are not on the real axis and we require

$$\left| \frac{2 - C_2 \pm j\sqrt{4C_1 - C_2^2}}{2} \right| < 1$$

which is equivalent to

$$(2 - C_2)^2 + (4C_1 - C_2^2) < 4$$

reducing to

$$C_1 < C_2$$

The solution is shown as the shaded region in Figure 5.

Combining the two graphs, we find the overall solution for a stable second-order DPLL in Figure 6. Mathematically, the DPLL parameters C_1 and C_2 must satisfy

$$C_1 > 0, C_1 > 2C_2 - 4, C_1 < C_2$$

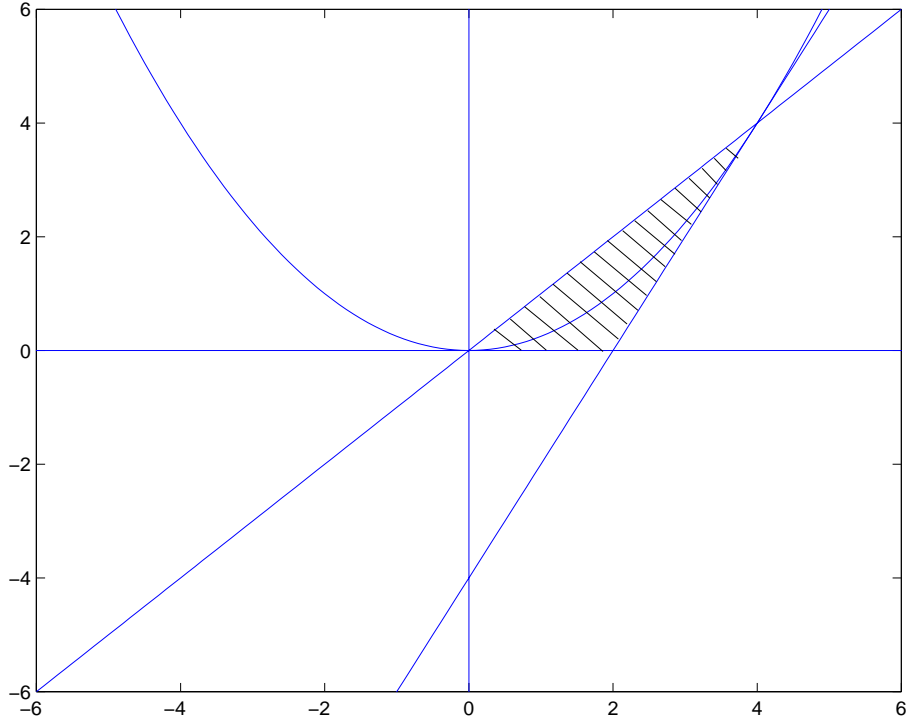


Figure 6: Second-order DPLL solution region. Horizontal axis is C_2 . Vertical axis is C_1 .

Referring to (5), we can also express everything in terms of natural frequency and damping factor. Here since C_1 and C_2 are both positive, ω_n and η must be of the same sign. It doesn't matter whether both of them are negative or positive because the transfer function only depends on ω_n^2 and $2\eta\omega_n$. We require both of them to be positive.

For the first case, we have $\eta > 1$, we have the stable condition

$$\begin{aligned}\omega_n &> 0 \\ \eta &> 1 \\ \eta\omega_n &< 2 \\ \omega_n^2 + 4 &> 4\eta\omega_n\end{aligned}$$

From the third inequality, we have

$$\omega_n < 2/\eta < 2$$

From the last inequality, we have

$$\eta\omega_n < \frac{\omega_n^2}{4} + 1$$

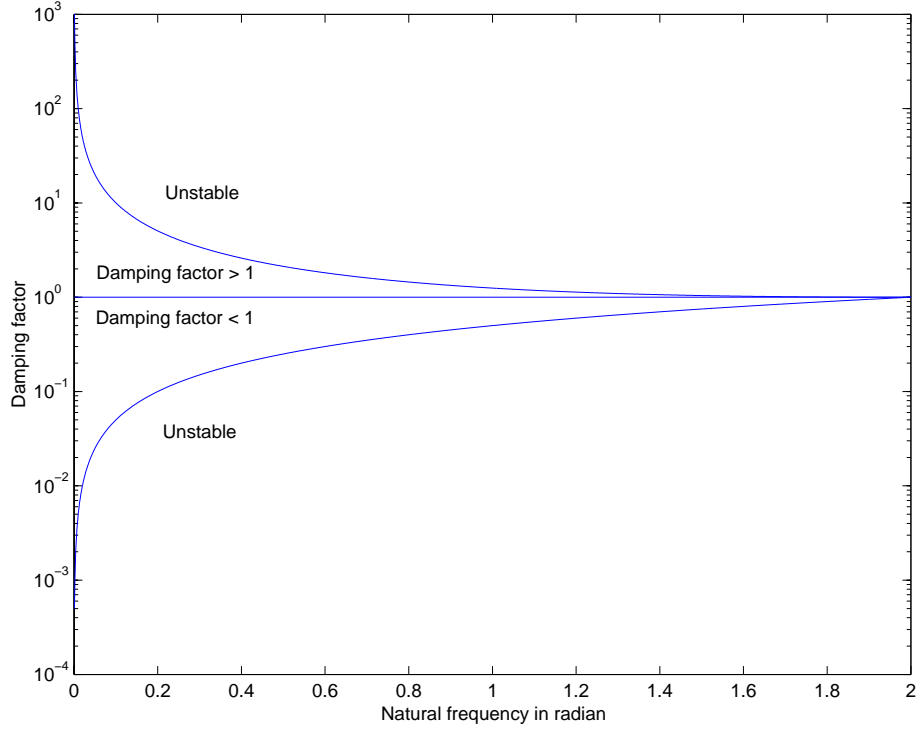


Figure 7: Second-order DPLL solution region in terms of natural frequency ω_n and damping factor η .

which combined with the condition $\omega_n < 2$ supercedes the second inequality above. We can thus rewrite the condition as

$$\begin{aligned} \eta &> 1 \\ 0 &< \omega_n < 2 \\ \eta\omega_n &< \frac{\omega_n^2}{4} + 1 \end{aligned} \quad (6)$$

For the second case, we have the stable condition

$$\begin{aligned} \eta &\leq 1 \\ 0 &< \omega_n < 2\eta \end{aligned} \quad (7)$$

The solution region in terms of natural frequency and damping factor is shown in Figure 7.

References

- [1] Edward A. Lee and David G. Messerschmit. *Digital Communications*. Kluwer Academic Publishers, second edition, 1994.
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- [3] Y. R. Shayan and T. Le-Ngoc. All digital phase-locked loop: concepts, design and application. *IEEE Proc.*, 136(1), February 1989.
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